# ORIGINAL PAPER

# Theoretical study of photophysical properties of 1,4-dihydropyrrolo [3,2-b]pyrrole-cored branched molecules with thienylenevinylene arms toward broad absorption spectra for solar cells

Shanshan Tang • Binbin Tang • Dadong Liang • Guang Chen • Ruifa Jin

Received: 17 April 2013 / Accepted: 9 June 2013 / Published online: 3 July 2013 © Springer-Verlag Berlin Heidelberg 2013

Abstract A series of oligo(thienylenevinylene) derivatives with 1,4-dihydropyrrolo[3,2-b]pyrrole as core has been investigated at the PBE0/6-31G(d) and the TD-PBE0/6-31+ G(d,p) levels to design materials with high performances such as broad absorption spectra and higher balance transfer property. The results show that position and amount of arm affect the electronic density contours of frontier molecular orbitals significantly. The molecule with four arms owns the narrowest energy gap and the largest maximum absorption wavelength, and the molecule with two arms in positions *a* and *c* has the broadest absorption region among the designed molecules. Calculated reorganization energies of the designed molecules indicate that the molecules with two

**Electronic supplementary material** The online version of this article (doi:10.1007/s00894-013-1920-y) contains supplementary material, which is available to authorized users.

S. Tang · D. Liang (⊠) College of Resource and Environmental Science, Jilin Agricultural University, Changchun 130118, China e-mail: dadongliang@126.com

#### B. Tang

School of Materials Science and Engineerning, Southwest University, Chongqing 400715, China

# G. Chen (🖂)

College of Life Sciences, Jilin Agricultural University, Changchun, Jilin 130118, China e-mail: chg61@163.com

### R. Jin

College of Chemistry and Chemical Engineering, Chifeng University, Chifeng 024000, China

arms can be good potential ambipolar transport materials under proper operating conditions.

**Keywords** Charge transport property · Electronic and optical properties · Organic solar cells · Thienylenevinylene derivatives

# Introduction

So far, a lot of effort has been devoted to the development of high-performance organic solar cell (OSC) materials [1-11]. For example, some studies indicate that branched oligothiophene and oligo(thienylenevinylene) derivatives can show excellent properties for OSC materials [12–14]. A theoretical study reported that 1,4-dihydropyrrolo[3,2-b]pyrrole-cored molecules were potential candidates for OSC materials [15, 16]. However, there still exist some challenges in the development of OSC materials. One of them is that, in present OSC devices, the photocurrent is limited because the overlap between the spectra of the sun and the absorption spectra of present organic materials for OSCs is not good [17]. The poor spectral overlap leads to the loss of the energy in the sunlight. The photons with energies below about 2.0 eV cannot be absorbed by the organic materials. Therefore, the spectral overlap is necessary to be increased, namely, the energy gap of  $\pi$ -conjugated molecule is necessary to be reduced. Understanding the relationship between molecular structure and charge transport property of material is a key factor for designing good candidates for OSC devices. The lower the reorganization energy, the higher the charge transfer rate [15, 16]. In this study, the charge transfer reaction is the

self-exchange reaction. In other words, the free energy difference ( $\Delta G^{\circ}$ ) between the initial and final states is approximately assumed as zero in the transfer process. As a result, the charge transfer rate can be described by Marcus theory [18] via the following equation:

$$K = V^2 / \hbar (\pi / \lambda k_{\rm b} T)^{1/2} \exp(-\lambda / 4k_{\rm b} T)$$
(1)

where T is the temperature,  $k_{\rm b}$  means the Boltzmann constant,  $\lambda$  represents the reorganization energy due to geometric relaxation accompanying charge transfer, and V is the electronic coupling matrix element (transfer integral) between the two adjacent species dictated largely by orbital overlap. It can be seen from Eq. (1) that there are two major parameters that determine the charge transfer rates: (i) V, which needs to be maximized, and (ii)  $\lambda$ , which needs to be small for significant transport. In order to investigate V, crystal data in general is required [19, 20]. However, the designed molecules may be non-crystal, and the electronic coupling matrix element value is very limited [21-24]. Though charge transfer integrals of amorphous dendrimers can be estimated [25-27], we cannot predict them for our designed molecules by similar methods due to the lack of available approaches. Thus, we focused on their reorganization energies to investigate their charge transport properties. Generally, the reorganization energy is determined by the fast change of the molecular geometry when a charge is added or removed from a molecule (the inner reorganization energy,  $\lambda_{int}$ ) and represents variations in the surrounding medium due to the polarization effects (the external reorganization energy,  $\lambda_{ext}$ ) [28-31]. Our designed molecules may be used as the donors of solar cells in solid films, because the dielectric constant of medium for molecule is low [21, 22]. The computed values of the external reorganization energy in pure organic condensed phases are not only small but also much smaller than their internal counterparts [32-34]. Moreover, previous works reported that there is a clear correlation between  $\lambda_{int}$  and charge transfer rate [25, 35, 36]. Therefore, we only paid attention to the discussion of the  $\lambda_{int}$  (reflects the geometric changes in the molecules when going from the neutral to the ionized state and vice versa) of the isolated active organic  $\pi$ -conjugated systems due to ignoring any environmental relaxation and changes in this paper. Hence, the electron reorganization energy ( $\lambda_e$ ) and hole reorganization energy ( $\lambda_{\rm h}$ ) values can be calculated by Eqs. (2) and (3) [37]:

$$\lambda_{\rm e} = \left[ E_0^- - E_- \right] + \left[ E_-^0 - E_0 \right] \tag{2}$$

$$\lambda_{\rm h} = \left[ E_0^+ - E_+ \right] + \left[ E_+^0 - E_0 \right]. \tag{3}$$

Where  $E_0^+$  ( $E_0^-$ ) is the energy of the cation (anion) calculated with the optimized structure of the neutral molecule. Similarly,  $E_+$  ( $E_-$ ) is the energy of the cation (anion) calculated with the optimized cationic (anionic) structure,  $E_{+}^{0}$  ( $E_{-}^{0}$ ) is the energy of the neutral molecule calculated at the cationic (anionic) state. Finally,  $E_{0}$  is the energy of the neutral molecule at the ground state.

The ionization potential (*IP*) and electron affinity (*EA*) of the molecules were calculated as described in the formulas below:

$$IP = E(\mathbf{M}^+) - E(\mathbf{M}_{\text{neutral}}) \tag{4}$$

$$EA = E(\mathbf{M}_{\text{neutral}}) - E(\mathbf{M}^{-})$$
(5)

where  $E(M_{neutral})$  is the total energy of the neutral form at the optimized geometry,  $E(M^+)$  and  $E(M^-)$  are the total energies of the cationic and the anionic forms of the molecules, respectively.

In this contribution, we designed three groups of oligo(thienylenevinylene) derivatives with 1,4-dihydropyrrolo [3,2-b]pyrrole as the core and thienylenevinylene as the arms (see Scheme 1). The first group is the derivatives with two arms: a-b, a-c, and a-d are di-substitutions with thienylenevinylene arms in positions a and b, a and c, and a and d, respectively. The second group is the derivatives with three arms: a-b-c and a-b-d correspond to tri-substitutions with thienylenevinylene arms in positions a, b, and c and a, b, and d, respectively. The third group is molecules with four arms: a-b-c-d refers to tetra-substitutions of thienylenevinylene arms in positions a, b, c, and d. The highest occupied molecular orbital (HOMO) energies, the lowest unoccupied molecular orbital (LUMO) energies, the HOMO-LUMO gaps  $(E_{\sigma})$ , and the absorption spectra were predicted. The electronic properties including natural bond orbital (NBO) charge, ionization potential (IP), electron affinity (EA), and reorganization energy ( $\lambda$ ) were investigated. In addition, the correlation between these properties and the molecular structures was discussed.

# **Computational details**

The quantum chemistry calculations of ground state were carried out with the aid of Gaussian 03 package [38]. The PBE0 method with the 6-31G(d) basis set was used in all the geometry optimization including neutral, cationic, and anionic molecules. A large number of previous works have suggested that the PBE0 method notably is adapted to the sulfur-bearing molecules [4–7, 15, 16, 39–44]. Frequency



Core: 1,4-dihydropyrrolo[3,2-b]pyrrole

Arm: Thienylenevinylene

Scheme 1 Chemical structures of core and arm of investigated molecules. Core: 1,4-dihydropyrrolo[3,2-b]pyrrole Arm: Thienylenevinylene

calculations using the same methods as those for the geometry optimizations were performed for the investigated molecules. All real frequencies have confirmed the presence of a local minimum. The absorption spectra of the designed molecules were investigated by employing the TD-PBE0/6-31+G(d,p) method based on the optimized geometries obtained at the PBE0/6-31G(d) level. For comparing with the interested results reported previously [45, 46], the reorganization energies for electron ( $\lambda_e$ ) and hole ( $\lambda_h$ ) of the molecules were predicted from the single point energy at the B3LYP/6-31G(d,p) level. The natural bond orbital (NBO) analysis was used to map the charge distribution of the investigated molecules. Gaussian 09 package [47] was employed to optimize the geometry of molecules in excited state and calculate the NBO charge of ground and excited states.

# **Results and discussion**

# Frontier molecular orbitals

In order to gain insight into the influence of the optical and electronic properties, the electronic density contours of

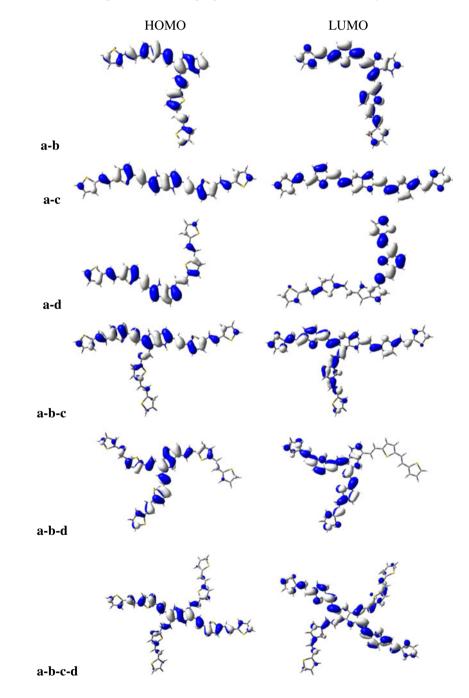
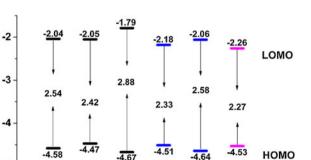


Fig. 1 Electronic density contours of the frontier molecular orbital for investigated molecules in ground state frontier molecular orbitals (FMOs) for the designed molecules in ground and excited states are investigated, and their sketches are plotted in Fig. 1 and Fig. S1, respectively. The evaluations of FMO energies for designed molecules are shown in Fig. 2.

As shown in Fig. 1, for the molecules with two arms, the electronic density contours of FMOs of a-b and a-c are mainly delocalized on the whole molecules. The electronic density contours of HOMO of **a-d** is mainly spread over the core, one arm in position a, and half arm in position d. The electronic density contours of LUMO of a-d is mainly delocalized on one arm in position d, and with minor contributions of the core and the arm in position a. For the three arms molecules, the electronic density contours of FMO of **a-b-c** are mainly distributed on the core, two arms in positions a and c, and half arm in position b. The electronic density contours of HOMO of a-b-d is mainly centralized on the core, two arms in positions a and b, and half arm in position d. The electronic density contours of LUMO of a-b**d** is mainly delocalized on the core and two arms in positions a and b. The electronic density contours of FMO of **a-b-c-d** are mainly spread over the core, two arm in positions a and c, and half arm in positions b and d. These results reveal that the position and amount of arm affect the distributions of the electronic density contours of FMO significantly. Except a-d displays a long-range electronic density reorganization during the excitation, the charge transfers of other molecules are rather localized, somehow being an undesired feature regarding the target electronic properties of the designed materials. That is to say, **a-d** will show better electron property.

The FMO energies are very important because they not only closely relate to the spectral properties but also to the reactivity of a compound. As shown in Fig. 2, for the two arms molecules, the HOMO energies ( $E_{HOMO}$ ) are in the order of  $\mathbf{a} - \mathbf{c} > \mathbf{a} - \mathbf{b} > \mathbf{a} - \mathbf{d}$ , and their LUMO energies ( $E_{LUMO}$ ) are in the order of  $\mathbf{a} - \mathbf{d} > \mathbf{a} - \mathbf{b} > \mathbf{a} - \mathbf{c}$ . Thus, their  $E_g$  values are in the order of  $\mathbf{a} - \mathbf{d} > \mathbf{a} - \mathbf{b} > \mathbf{a} - \mathbf{c}$ . Thus, their  $E_g$  values are in the order of  $\mathbf{a} - \mathbf{d} > \mathbf{a} - \mathbf{b} > \mathbf{a} - \mathbf{c}$ . These results display that for the two arms molecules, the arms in a and c positions result in the narrowest HOMO-LUMO gap. For the three arms



a-b-c

a-b-d a-b-c-d

Fig. 2 Evaluation of the calculated FMOs energies for investigated molecules at the PBE0/6-31G(d) level

a-d

-5

a-b

a-c

Energy (eV)

molecules, the  $E_{\text{HOMO}}$ ,  $E_{\text{LUMO}}$ , and  $E_{\text{g}}$  of **a-b-c** are higher, lower, and narrower than that of **a-b-d**, respectively. It indicates that the arm in position c leads to smaller  $E_{\text{g}}$  value than the arm in position d. The values of  $E_{\text{HOMO}}$ ,  $E_{\text{LUMO}}$ , and  $E_{\text{g}}$  of **a-b-c** are larger, smaller, and smaller than that of **a-b**, respectively. It shows that increasing the arm in position cresults in smaller  $E_{\text{g}}$  value. **a-b-c-d** owns the lowest  $E_{\text{LUMO}}$ and the narrowest  $E_{\text{g}}$  among the designed molecules. It should lead to the longest wavelength of absorption spectrum.

# Absorption spectra

The values of longest wavelength of absorption spectrum  $(\lambda_{\text{max}})$ , corresponding oscillator strength (*f*), vertical excitation energy ( $E_v$ ), and absorption region (*R* denotes for the difference of the longest and shortest wavelength values with oscillator strength larger than 0.01 considering the first twenty excited states, see Table S1, the Supporting information) of the designed molecules are listed in Table 1. The simulated absorption spectra of the investigated molecules are shown in Fig. 3.

For the molecules with two arms, the  $\lambda_{max}$ , corresponding f, and R values are all in the order of  $\mathbf{a} - \mathbf{c} > \mathbf{a} - \mathbf{b} > \mathbf{a} - \mathbf{d}$ , which is in good agreement with the corresponding reverse order of their  $E_{\rm v}$  values. It reveals that the arms in positions a and c may result in large  $\lambda_{\max}$  value and broad absorption region. For the three arms molecules, the  $\lambda_{\rm max}$ , corresponding f, and R values are all in the order of **a-b-c** > **a-b-d**, and the  $E_v$  value of **a-b-c** is smaller than that of **a-b-d**. It shows that the arm in position c may lead to larger  $\lambda_{\text{max}}$ , f, and R values than the arm in position d, respectively. The  $\lambda_{\text{max}}$ , corresponding f, and R values of a-b-c are larger than that of a-b, which displays that increasing the arm in position c may increase the values of these properties. **a-b-c-d** has smaller  $E_{\rm v}$  and larger  $\lambda_{\rm max}$ values than that of a-b-c, which indicates that increasing the arm in position d can increase the  $\lambda_{max}$  value. In sum, the position and amount of the arm significantly influences optical properties of these designed molecules. As a

Table 1 Optical properties computed at the TD-PBE0/6-31+G(d,p)// PBE0/6-31G(d) level

_	$E_{\rm v}({\rm eV})$	$\lambda_{\max}$ (nm)	f	<i>R</i> (nm)	
a-b	2.15	577.77	1.61	298.51	
a-c	2.07	599.40	3.33	314.80	
a-d	2.43	510.20	1.27	232.59	
a-b-c	1.98	627.15	2.11	312.61	
a-b-d	2.12	584.21	0.91	272.09	
a-b-c-d	1.94	640.68	2.05	286.89	

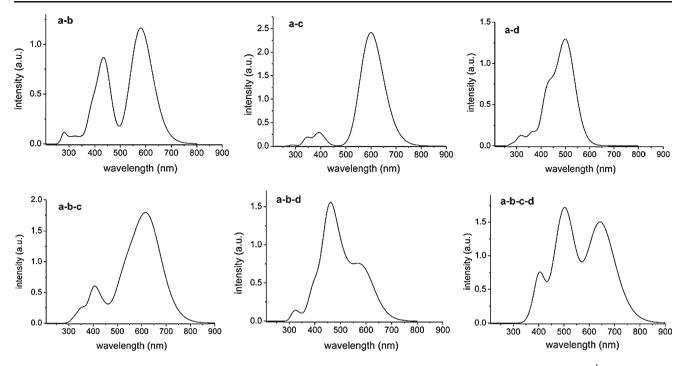


Fig. 3 Simulated absorption spectra of investigated molecules. The value of the full width at half maximum (fwhm) is 3000 cm<sup>-1</sup>

result, **a-b** has the largest f and R values, and **a-b-c-d** has the largest  $\lambda_{\text{max}}$  value.

### Charge transport properties

To compare the charge transport properties of the investigated molecules, we calculated the reorganization energies associated with different geometries of two states (anion and cation) based on the Eqs. (2) and (3). The results are summarized in Table 2. Herein, the reorganization energy is just the internal reorganization energy of the isolated active organic  $\pi$ -conjugated systems due to ignoring any environmental relaxation and changes. In order to study the response of the molecule to the formation of a hole, or to the addition of an electron, *IP* and *EA*, both vertical (v: at the geometry of the neutral molecule) and adiabatic (a: at the optimized structures of both the neutral and charged molecule) Eqs. (4) and

**Table 2** Calculated molecular ionization potential ( $IP_a$  and  $IP_v$ ), electron affinity ( $EA_a$  and  $EA_v$ ), reorganization energies ( $\lambda_e$  and  $\lambda_h$ ) computed at the B3LYP/6-31G(d,p)//PBE0/6-31G(d) level (in eV)

	IPa	$I\!P_{\rm v}$	EA <sub>a</sub>	$EA_{\rm v}$	$\lambda_{\rm e}$	$\lambda_{ m h}$
a-b	5.230	5.335	1.308	1.124	0.269	0.223
a-c	5.103	5.202	1.365	1.202	0.224	0.214
a-d	5.350	5.427	1.034	0.913	0.143	0.186
a-b-c	5.043	5.128	1.606	1.419	0.259	0.197
a-b-d	5.194	5.265	1.419	1.251	0.245	0.171
a-b-c-d	4.990	5.066	1.740	1.570	0.223	0.169

(5) were computed. The calculated results are listed in Table 2. The natural bond orbital (NBO) analysis was supplied to map the charge distribution of the models and provide the evidence of charge transfer (see Fig. S2 and Tables S2-4, the Supporting information).

It is well known the lower the IP of the hole-transport layer (HTL), the easier the entrance of holes from HTL to indium tin; the lower the EA of the electron-transport layer (ETL), the easier the entrance of electrons from ETL to cathode [37]. From the results in Table 2, one can see that the values of  $IP_{a}$ and  $IP_v$  are both in the order of  $\mathbf{a} - \mathbf{d} > \mathbf{a} - \mathbf{b} > \mathbf{a} - \mathbf{c}$ , and the values of  $EA_{a}$  and  $EA_{v}$  are both in the order of  $\mathbf{a} - \mathbf{d} < \mathbf{a} - \mathbf{b} < \mathbf{a} - \mathbf{c}$ . These results show that the orders of  $IP_a$  and  $IP_v$  values are nearly reversed with their orders of  $EA_a$  and  $EA_v$  values. On one hand, the  $IP_{\rm a}$  and  $IP_{\rm v}$  values of **a-b-c** are smaller than that of **ab-d**. On the other hand, the  $EA_a$  and  $EA_v$  values of **a-b-c** are larger than that of **a-b-d**, respectively. The  $IP_a$  and  $IP_v$  values of **a-b-c** are smaller than that of **a-b**, respectively. It reveals that increasing the arm in position c can decease the  $IP_a$  and  $IP_{\rm v}$  values. Among these molecules, **a-d** has the smallest  $EA_{\rm a}$ and  $EA_v$  values, and **a-b-c-d** owns the smallest  $IP_a$  and  $IP_v$ values. It displays that **a-d** is the most proper one as ETL material, and **a-b-c-d** is the most proper one as HTL material.

For one thing, the  $\lambda_e$  values of designed molecules (0.143– 0.269 eV) are smaller than that of tris(8-hydroxyquinolinato) aluminum(III) (Alq3) ( $\lambda_e$ =0.276 eV) which is a typical electron transport material [45]. It indicates that their electron transfer rates are higher than that of Alq3. On the other hand, their  $\lambda_h$  values (0.169–0.223 eV) are smaller than that of N,N'diphenyl-N,N'-bis(3-methlphenyl)-(1,1'-biphenyl)-4,4'- diamine (TPD) ( $\lambda_{\rm h}$ =0.290 eV) which is a typical hole transport material [46]. It implies that their hole transfer rates are higher than that of TPD. Moreover, the differences between  $\lambda_e$ and  $\lambda_h$  values for **a-b**, **a-c**, and **a-d** are 0.046, 0.010, and 0.043 eV, respectively. It reveals that they have better equilibrium properties for hole- and electron-transport. As a result, these molecules are potential ambipolar charge transport materials under proper operating conditions. The values of  $\lambda_e$  and  $\lambda_h$  are both in the order of **a-b** > **a-c** > a-d, which shows that the arms in positions a and d may lead to better charge transfer rate. The values of  $\lambda_e$  and  $\lambda_h$ of **a-b-c** are larger than that of **a-b-d**, respectively. Among these designed molecules, **a-d** has the smallest  $\lambda_e$  value, and **a-b-c-d** owns the smallest  $\lambda_h$  value. These designed molecules are potential candidates as hole transfer materials toward solar cells expect **a-d**, because their  $\lambda_{\rm h}$  values are smaller than their  $\lambda_e$  values, respectively.

From Figs. 1 and S1 (Supporting information), one can find that the electronic density contours of FMOs for molecules in excited state are similar in ground state, respectively. These results show that the donor and accepter are the whole molecules for **a-b** and **a-c**. The donor of **a-d** is the core, one arm in position a, and half arm in position d, and its acceptor is the arm in position d. The donor and acceptor of **a-b-c** are the core, two arms in positions a and c, and half arm in position b. The donor of **a-b-d** is the core, two arms in positions a and b, and half arm in position d, and its acceptor is the core and two arms in positions a and b. The donor and acceptor of **a-b-c-d** are the core, two arms in positions a and c, and half arm in positions b and d. Figure S2 and Tables S2-4 (Supporting information) suggest that there is charge transfer during the excitation. It agrees with the results obtained in electronic density contours of FMOs for molecules.

## Conclusions

In the present work, we report a theoretical investigation of predicting the effects of position and amount of arm on optical and electronic properties for a series of 1,4dihydropyrrolo[3,2-b]pyrrole-cored branched molecules with thienylenevinylene arms. The following conclusions can be drawn. (1) The position and amount of arm affect the distributions of the FMO patterns significantly. (2) a-b-c-d owns the lowest  $E_{\text{LUMO}}$  and the narrowest  $E_{\text{g}}$ , and increasing the arm in position c results in smaller  $E_{g}$  value. (3) **a-b** has the largest f and R values, and **a-b-c-d** has the largest  $\lambda_{\text{max}}$  value. (4) **a-d** is the most proper one as ETL material, and **a-b-c-d** is the most proper one as HTL material. (5) a-d has the smallest  $\lambda_{\rm e}$  value, and **a-b-c-d** owns the smallest  $\lambda_{\rm h}$  value. These results can be helpful for further theoretical and experimental study of oligo(thienylenevinylene) derivatives toward the application as OSC materials.

Acknowledgments Financial support from the Science Foundation for Young Teachers of Jilin Agricultural University (No. 201219) is gratefully acknowledged.

#### References

- Kim JY, Lee K, Coates NE, Moses D, Nguyen TQ, Dante M, Heeger AJ (2007) Science 317:222–225
- Sang-aroon W, Laopha S, Chaiamornnugool P, Tontapha S, Saekow S, Amornkitbamrung V (2013) J Mol Model 19:1407–1415
- Zhang WB, Tu YF, Sun HJ, Yue K, Gong X, Cheng Stephen ZD (2012) Sci China Chem 55:749–7540
- 4. Tang SS, Zhang JP (2011) Inter J Quantum Chem 111:2089-2098
- 5. Tang SS, Zhang JP (2011) J Phys Chem A 115:5184-5191
- Liang DD, Liu JB, Kang LJ, Jin RF, Tang SS (2012) Mol Phys 110:369–375
- Liang DD, Tang SS, Liu JB, Liu JH, Kang LJ (2009) J Mol Struct THEOCHEM 908:102–106
- 8. Cheng YJ, Yang SH, Hsu CS (2009) Chem Rev 109:5868-5923
- Gutiérrez-Pérez RM, Flores-Holguín N, Glossmann-Mitnik D, Rodriguez-Valdez LM (2011) J Mol Model 17:1963–1972
- 10. López AM, Mateo-Alonsobe A, Prato M (2011) J Mater Chem 21:1305–1318
- Zhang CR, Liu L, Zhe JW, Jin NZ, Yuan LH, Chen YH, Wei ZQ, Wu Y-Z, Liu ZJ, Chen HS (2013) J Mol Model 19:1553–1563
- Fichou D (ed) (1999) Handbook of oligo- and polythiophenes. Wiley-VCH, Weinheim
- Vélez JH, Gutiérrez-Oliva S, Díaz FR, Angelica del Valle M, Toro-Labbé A, Bernède JC, East GA (2011) J Mol Model 17:81–88
- Müllen K, Wegner G (eds) (1998) Electronic materials: The oligomer approach. Press: Wiley-VCH, Weinheim
- 15. Tang SS, Zhang JP (2012) J Comput Chem 33:1353-1363
- Hu B, Yao C, Wang QW, Zhang H, Yu JK (2012) Sci China Chem 55:1364–1369
- 17. Winder C, Sariciftci NS (2004) J Mater Chem 14:1077-1086
- 18. Marcus RA (1993) Rev Mod Phys 65:599-610
- 19. Marcus RA (1964) Theo Annu Rev Phys Chem 15:155-196
- 20. Hush NS (1958) J Chem Phys 28:962
- Yang XD, Wang LJ, Wang CL, Long W, Shuai ZG (2008) Chem Mater 20:3205–3211
- 22. Wang CL, Wang FH, Yang XD, Li QK, Shuai ZG (2008) Org Electron 9:635–640
- 23. Troisi A, Orlandi G (2006) J Phys Chem A 110:4065-4070
- Yin SW, Yi YP, Li QX, Yu G, Liu YQ, Shuai ZG (2006) J Phys Chem A 110:7138–7143
- Köse ME, Mitchell WJ, Kopidakis N, Chang CH, Shaheen SE, Kim K, Rumbles G (2007) J Am Chem Soc 129:14257–14270
- Nelsen SF, Trieber DA, Ismagilov RF, Teki Y (2001) J Am Chem Soc 123:5684–5694
- 27. Nelsen SF, Blomgren F (2001) J Org Chem 66:6551-6559
- Käse ME, Long H, Kim K, Graf P, Ginley D (2010) J Phys Chem A 114:4388–4393
- Köse ME, Graf P, Kopidakis N, Shaheen SE, Kim K, Rumbles G (2009) Chemphyschem 10:3285–3294
- Lemaur V, Steel M, Beljonne D, Brédas JL, Cornil J (2005) J Am Chem Soc 127:6077–6086
- 31. Cheung DL, Troisi A (2010) J Phys Chem C 114:20479-20488
- 32. Norton JE, Brédas JL (2008) J Am Chem Soc 130:12377-12384
- Martinelli NG, Idé J, Sánchez-Carrera RS, Coropceanu V, Brédas JL, Ducasse L, Castet F, Cornil J, Beljonne D (2010) J Phys Chem C 114:20678–20685
- 34. Mc Mahon DP, Troisi A (2010) J Phys Chem Lett 1:941-946

- 35. Sakanoue K, Motoda M, Sugimoto M, Sakaki S (1999) J Phys Chem A 103:5551–5556
- 36. Hutchison GR, Ratner MA, Marks TJ (2005) J Am Chem Soc 127:2339–2350
- Zou LY, Ren AM, Feng JK, Liu YL, Ran XQ, Sun CC (2008) J Phys Chem A 112:12172–12178
- Frisch MJ et al. (2003) Gaussian 03, Revision B.03. Gaussian, Inc, Pittsburgh
- 39. Gahungu G, Zhang B, Zhang JP (2007) J Phys Chem C 111:4838-4846
- 40. Hu B, Yao C, Huang XR (2012) Spectrosc Lett 45:17-21
- 41. Jacquemin D, Perpete EA (2006) Chem Phys Lett 429:147–152

- 42. Perpete EA, Preat J, Andre JM, Jacquemin D (2006) J Phys Chem A 110:5629–5635
- 43. Jacquemin D, Wathelet V, Perpete EA (2006) J Phys Chem A 110:9145–9152
- 44. Jacquemin D, Preat J, Wathelet V, Fontaine M, Perpete EA (2006) J Am Chem Soc 128:2072–2083
- 45. Lin BC, Cheng CP, You ZQ, Hsu CP (2005) J Am Chem Soc 127:66–67
- Gruhn NE, da Silva Filho DA, Bill TG, Malagoli M, Coropceanu V, Kahn A, Brédas JL (2002) J Am Chem Soc 124:7918–7919
- 47. Frisch MJ et al. (2009) Gaussian 09. Gaussian, Wallingford